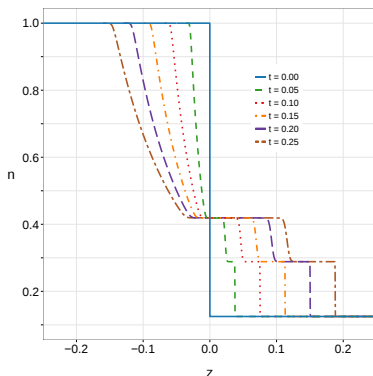


# Simulating relativistic flows with a Lattice Boltzmann model based on quadratures



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- 1 Relativistic Boltzmann equation
- 2 Our Lattice Boltzmann model
- 3 Results: The Riemann problem

# Boltzmann equation

- Tool for investigating fluids at a **mesoscopic** level
- Different levels
  - Macroscopic: continuous fields (hydrodynamics)
  - Mesoscopic: distribution function (*Boltzmann equation*)
  - Microscopic: individual molecules (molecular dynamics)
- Provides a good description of fluids from the **hydrodynamic** to the **ballistic** regime

# Relativistic Boltzmann equation

- Boltzmann equation for 1D flow in flat spacetime:

$$(\partial_t + v^z \partial_z) f = -\frac{u^0 - v^z u^z}{\tau} (f - f^{\text{eq}}),$$

$$v^z \equiv \xi = \frac{p^z}{p^0}, \quad p^0 = p \equiv |\vec{p}|^2. \quad (1)$$

- The relativistic Maxwell-Jüttner (M-J) distribution function for massless particles:

$$f^{\text{eq}} = \frac{n}{8\pi T^3} \exp\left(\frac{p^{\hat{\alpha}} u_{\hat{\alpha}}}{T}\right). \quad (2)$$

# Quadrature method + finite difference

- We project the distribution functions onto a set of orthogonal polynomials, and truncate the series

$$f^{eq}(x, \vec{p}, t) = e^{-p} \sum_{k=0}^{N_p-1} \sum_{n=0}^{N_v-1} a_{(nk)}^{i_1 \dots i_n}(x, t) P_{i_1 \dots i_n}^{(n)}(\vec{v}) L_j^{(1)}(p) \quad (3)$$

- Use Gauss-quadrature method to write integrals as sums<sup>1,2</sup>

$$\begin{aligned} \mathcal{Q}(p, \theta) &= \int_0^{2\pi} d\phi f^{eq} A_n(p, \xi, \phi) = \sum_{i=0}^{Q_\phi-1} \frac{2\pi}{Q_\phi} f_i^{eq} A_n(p, \theta, \phi_i) \\ \mathcal{E}(p) &= \int_{-1}^1 d(\cos \theta) \mathcal{Q}(p, \theta) = \sum_{j=0}^{Q_\xi-1} w_j^\xi \mathcal{Q}(p, \theta_j) \\ \mathcal{M} &= \int_0^\infty dp p e^{-p} \mathcal{E}(p) = \sum_{k=0}^{Q_p-1} w_k^p \mathcal{E}(p_k) \end{aligned} \quad (4)$$

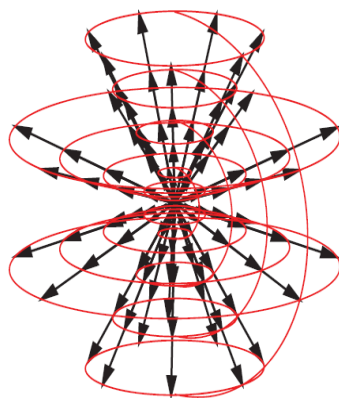
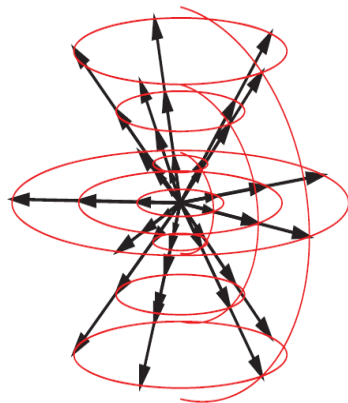
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<sup>1</sup>V.E. Ambrus, V. Sofonea, Phys. Rev. E **86**, 016708 (2012)

<sup>2</sup>P. Romatschke, M. Mendoza, and S. Succi, Phys. Rev. C **84**, 034903 (2012)

# Quadrature method

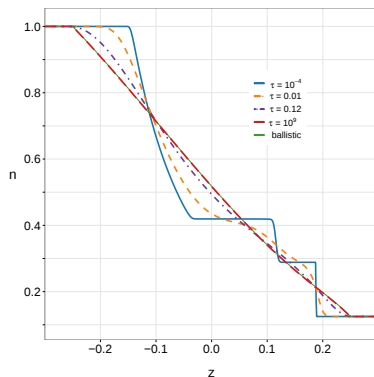
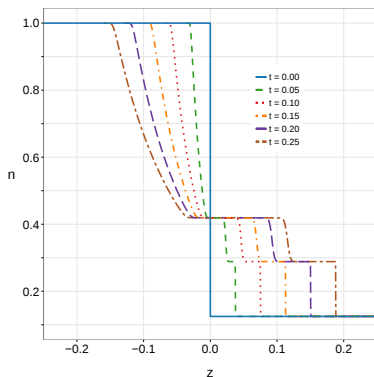
- Examples for sets of momenta:<sup>3</sup>



$$\vec{p} \rightarrow \{p_i, \xi_j, \phi_k\} \quad (5)$$

<sup>3</sup>V.E. Ambruş, V. Sofonea, Phys. Rev. E **86**, 016708 (2012)

# The Riemann problem

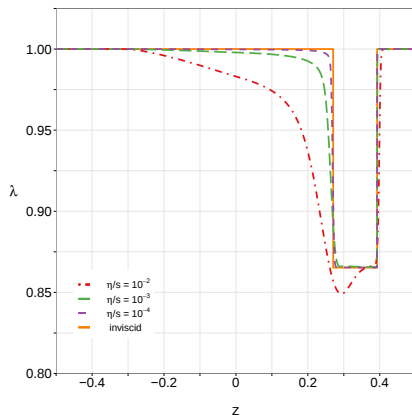
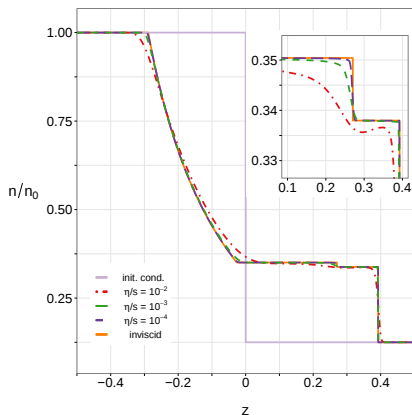


- Time evolution of the Sod shock tube problem<sup>4</sup>: formation of a shock-wave, contact discontinuity and rarefaction wave<sup>5</sup>
- As  $\tau$  increases, the features are smoothed, while finally in the ballistic regime they do no longer form.

<sup>4</sup>R. Blaga and V. Ambrus, arXiv:1608.03004 (2016)

<sup>5</sup>L. Rezzolla and Z. Olindo, *Relativistic hydrodynamics*, Oxford University Press, 2013.

# Inviscid regime ( $\tau \rightarrow 0$ ): density and fugacity

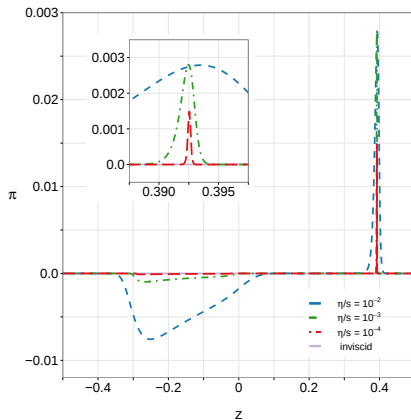
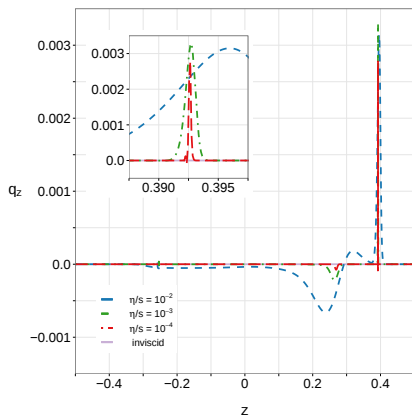


- density ( $n$ ) and fugacity ( $\lambda$ ) profiles for different relaxation times<sup>6</sup>

$$\frac{\eta}{s} \simeq \frac{1}{6} \tau T, \quad \lambda = e^{\mu_E/T} \equiv \frac{1}{d_G} \frac{n}{8\pi T^3}. \quad (6)$$

<sup>6</sup>Initial conditions taken from: I. Bouras et al. , Phys. Rev. Lett. **103** (2009).

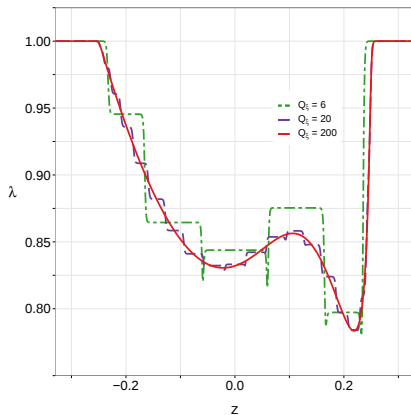
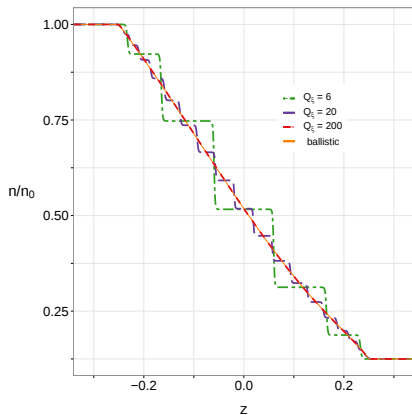
# Inviscid regime ( $\tau \rightarrow 0$ ): heat flow and shear pressure



- parallel component of the heat flux ( $q_z$ ) and shear pressure ( $\pi$ ), for decreasing values of the relaxation time

<sup>6</sup>I. Bouras et al. , Phys. Rev. Lett. **103** (2009).

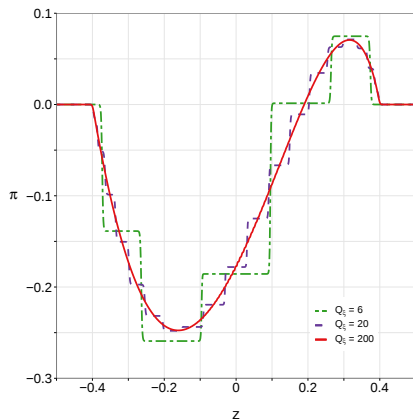
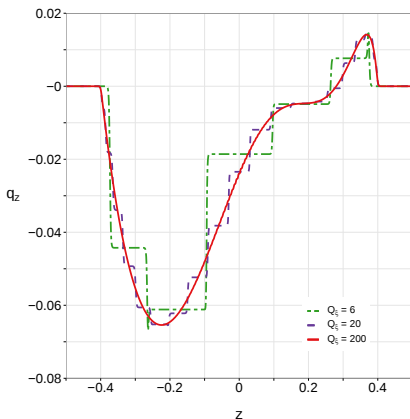
# Ballistic regime ( $\tau \rightarrow \infty$ ): density and fugacity



- density ( $n$ ) and fugacity ( $\lambda$ ) profiles in the ballistic regime, for increasing order of the polar quadrature

<sup>6</sup>I. Bouras et al. , Phys. Rev. Lett. **103** (2009).

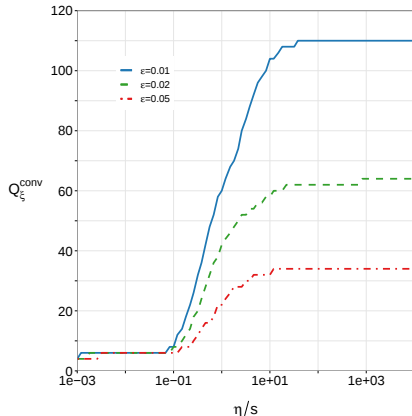
# Inviscid regime ( $\tau \rightarrow \infty$ ): heat flow and shear pressure



- parallel component of the heat flux ( $q_z$ ) and shear pressure ( $\pi$ ) profiles in the ballistic regime, for increasing order of the polar quadrature

<sup>6</sup>I. Bouras et al. , Phys. Rev. Lett. **103** (2009).

# Convergence test



- order of the polar quadrature necessary for convergence at a fixed accuracy

$$\varepsilon = \max_z \left[ \frac{n(z) - n_{\text{ref}}(z)}{\Delta n} \right]. \quad (7)$$

- nr. of velocities :

$$Q = Q_p * Q_\phi * Q_\xi$$

$$\text{inviscid} : Q = 3 * 2 * 4 = 24$$

$$\text{ballistic} : Q = 3 * 2 * 110 = 660. \quad (8)$$

# Summary

- we have developed a R-SLaB model based on Laguerre and Gauss quadratures
- numerical tests show good results across all regimes of the relaxation time (from hydro, through viscous, up to the ballistic regime)
- convergence test shows we need:
  - 24-36 velocities in the inviscid regime
  - $\sim 1000$  velocities in the ballistic regime
  - transition to higher orders around  $\tau \sim 0.01 - 0.1$ .

THANK YOU!

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