JUDD-OFELT ANALYSIS OF THE Er³⁺ IONS IN CaF₂ CRYSTALS

A. Lucaci, M. Stef, E. Preda, I. Nicoara

West University of Timisoara, Dept. of Physics, Timisoara

Abstract

Calcium fluoride is a well-known crystal that can be grown easily by Bridgman or Czochralski techniques; CaF_2 crystals has been used for long time in many optical components due to its exceptional transparency in the UV as well as in the IR. In this work a Judd-Ofelt investigation on the Er^{3+} ions in CaF_2 :ErF₃ single crystals with various ErF_3 concentrations have been performed. The intensity parameters are used to determine the emission probabilities of transitions and the branching ratios of the Er^{3+} transitions. The radiative life times of the excited states are also determined.

Keywords: Judd-Ofelt theory, intensity parameters, absorption spectrum, reduced matrix operator.

1. Introduction

The CaF₂ crystals have been used for long time in many optical components due to its exceptional transparency in the UV as well as in the IR spectral domain. Due to their optical properties, CaF₂ appears as a very attractive laser host which combines several advantages of calcium fluoride crystals and glasses. Another interest of CaF₂: Er^{+3} laser material, compared with other fluoride crystals resides in the possibility of the pumping with laser diodes and the choice of the excitation wavelength is not critical parameter because of the broad absorption bands. The optical properties of CaF₂ doped with trivalent rare-earth (RE) ions have been investigated by many authors [1,2,3]. In order to characterize the laser properties of these compounds we can use experimental or theoretical techniques.

The Judd-Ofelt model [4,5] permits to calculate some characteristics of the stimulated emission only from the absorption spectra [1,2,3]. In this paper we report a study about the influence of Er^{3+} ions concentration on the optical absorption spectra as well as the usage the Judd-Ofelt model in order to determine the intensity parameters Ω_i and the study the influence of the Er ³⁺ ion concentration on the Ω_i parameters.

2. Experimental procedure

Pure and ErF_3 – doped CaF_2 crystals have been grown in vacuum (10⁻¹Pa) in our crystal research laboratory using vertical Bridgman method [6]. A chemically pure grade (99,9%) ErF_3 reagent and suprapure grade Merck calcium fluoride were used as the starting materials. Transparent single crystals with various concentrations (0.15, 0.30 and 0.67 mol% ErF_3) in graphite crucible have been obtained. The rate of the crucible lowering was 4 mm/h.

The optical absorption spectra at room temperature have been obtained using a Shimadzu 1650PC spectrophotometer.



Fig. 1: Absorption spectra of various ErF₃ concentrations doped CaF₂ crystals.

3. Results and Discussions

In order to characterize the laser properties of laser materials we can use experimental or theoretical techniques. To investigate the stimulated emission some parameters must be known: the probability of the spontaneous emission $(A_{JJ'})$ and of the nonradiative transitions $(W_{JJ'})$, for example. Probability $A_{JJ'}$ is the sum of the probabilities of forced electric-dipole (ed) and magnetic-dipole (md) transitions; these parameters can be calculated if we known the line strength s_{JJ}^{ed} and s_{JJ}^{md} . In these expressions the Ω_t parameters appears, which are the intensity parameters describing the efficiency of the interaction between the RE³⁺ activator

ion and the surrounding crystal field. The Judd-Ofelt model [4, 5] permits to calculate these parameters from the investigation of the absorption spectra only.

The Judd-Ofelt model can be applied to the analysis of the absorption intensities of RE $(4f^N)$ transitions in the rare-earth-doped crystals, in order to determine the intensity parameters Ω_2 , Ω_4 , Ω_6 . Electric and magnetic dipole transitions between levels of the $4f^N$ configurations are perturbed by the static crystalline field.

The oscilator strength S_{calc}^{el} corresponding to the electric dipole transition from the ground level Ψ'_{r} is given by [1]:

$$S_{\text{calc}}^{\text{el}} = \sum_{t=2,4,6} \Omega_t \left| \left\langle \Psi_J \right| \left| U^{(t)} \right| \left| \Psi'_J \right\rangle \right|^2$$

where $U^{(t)}$ is a tensor operator of rank t (t = 2,4,6), and Ω_t are the so-called Judd-Ofelt parameters or intensity parameters.

Ten Er^{3+} absorption bands in the room-temperature absorption spectrum between 350 and 1000 nm (shown in Fig. 1), were chosen to determine Judd-Ofelt intensity parameters for Er^{3+} in CaF₂ for three different Er^{3+} concentrations: 0.15, 0.30 and 0.67 mol% ErF_3 . The peak wavelengths of the Er^{3+} bands are given in Table 1. The measured line strengths $\left(S_{JJ}^{\text{ed}}\right)_{\text{meas.}}$ of the chosen bands are determined using the following expression [1,3]:

$$\left(S_{JJ'}^{\text{ed}}\right)_{\text{meas.}} = \frac{1}{N_0} \int k(\lambda) d\lambda \frac{3 \cdot hc(2J+1)}{8\pi^3 e^2 \overline{\lambda}} \cdot \frac{9 \cdot n}{(n^2+2)^2}$$
(1)

where J and J' are the total angular momentum quantum numbers of the initial and final states, respectively, n is the refractive index, N_0 is the Er^{3+} ion concentration, $\overline{\lambda}$ is the mean wavelength of the specific absorption band and $\int k(\lambda) d\lambda$ is the integral of the absorption coefficient of the band associated with the transition $J \rightarrow J'$. The refractive indices of CaF2:Er³⁺ for all erbium concentrations was determined from Sellmeier's dispersion equation:

$$n(\overline{\lambda}_{JJ'}) = \sqrt{1 + \frac{B_1 \overline{\lambda}_{JJ'}^2}{\overline{\lambda}_{JJ'}^2 - C_1} + \frac{B_2 \overline{\lambda}_{JJ'}^2}{\overline{\lambda}_{JJ'}^2 - C_2} + \frac{B_3 \overline{\lambda}_{JJ'}^2}{\overline{\lambda}_{JJ'}^2 - C_3}}$$
(2)

where B_i and C_i (i = 1, 2, 3) are Sellmeier's coefficients: $B_1 = 5.675888 \times 10^{-1}$, $B_2 = 4.710914 \times 10^{-1}$, $B_3 = 3.8484723$, $C_1 = 2.52642999 \times 10^{-3}$, $C_2 = 1.00783328 \times 10^{-2}$ and $C_3 = 1.20055597 \times 10^3$.

Transition	$\overline{\lambda}$		$\begin{bmatrix} I I^{(2)} \end{bmatrix}^2$	$\begin{bmatrix} I I^{(4)} \end{bmatrix}^2$	$\begin{bmatrix} \mathbf{I} \\ \mathbf{I} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} \begin{bmatrix} \mathbf{I} \end{bmatrix} $	
from ${}^{4}I_{15/2}$	(nm)	n				
${}^{4}I_{11/2}$	977	1.429	0.0282	0.0003	0.3953	
$^{4}I_{9/2}$	805	1.430	0	0.1732	0.0099	
${}^{4}F_{9/2}$	655	1.432	0	0.5354	0.4619	
${}^{4}S_{3/2}$	545	1.435	0	0	0.2211	
$^{2}H(2)_{11/2}$	519	1.436	0.7125	0.4123	0.0925	
${}^{4}F_{7/2}$	487	1.437	0	0.1468	0.6266	
${}^{4}F_{5/2}$	447	1.439	0	0	0.2233	
$^{2}G(1)_{9/2}$	406	1.441	0	0.0190	0.2255	
${}^{4}G_{11/2}$	379	1.444	0.9181	0.5261	0.1171	
${}^{4}G_{9/2}$	360	1.445	0	0.2415	0.1234	

Table 1: Values of reduced matrix elements for the absorption bands of Er^{3+} in CaF_2 [1]

The measured line strengths were used to obtain the Judd-Ofelt parameters Ω_2 , Ω_4 and Ω_6 by solving a set of ten equations simultaneously for the corresponding transitions between *J* and *J*' manifolds in the following form:

$$\left(S_{JJ'}^{ed}\right)_{calc.} = \sum_{t=2,4,6} \Omega_t \left| \left\langle 4f^n \alpha[L,S]J \right| \left| U^{(t)} \right| \left| 4f^n \alpha'[L',S']J' \right\rangle \right|^2$$
(3)

where $\langle ||U^{(t)}|| \rangle$ are the doubly reduced matrix elements of rank t (t = 2, 4, 6) between states characterized by the quantum numbers (S, L, J) and (S', L', J'). The values of the reduced matrix elements are given in Table 1 and the values of the measured and calculated absorption line strengths are given in Table 2. The values of the J-O parameters calculated for three different erbium concentrations are give in Table 3. The J-O parameters can now be applied to calculate, using equation (3), the line strengths, corresponding to the transition from the upper manifold states, ${}^{4}I_{11/2}$, ${}^{4}I_{9/2}$, ${}^{4}F_{9/2}$, ${}^{4}S_{3/2}$, ${}^{2}H(2)_{11/2}$, ${}^{4}F_{7/2}$, ${}^{4}F_{5/2}$ and ${}^{2}G(1)_{9/2}$ to their corresponding lower-lying manifolds states. Using these line strengths, the radiative decay rates, $A_{JJ'}$, for electric dipole transitions between an excited states (J) and the lower-lying terminal manifolds (J') can be calculated using the expression [1]:

$$A_{\rm JJ'} = \frac{16\pi^3 e^2}{3h(2J+1)} \frac{n(n^2+2)^2}{9} \left(S_{\rm JJ'}^{\rm ed}\right)_{\rm calc.} \tag{4}$$

Concentration of Er ³⁺	0.15 1	mol%	0.30	mol%	0.67 mol%		
Transition	$\left(S_{\rm JJ'}^{\rm ed}\right)_{\rm meas.}$	$\left(S_{\rm JJ'}^{\rm ed} ight)_{ m calc.}$	$\left(S_{\rm JJ'}^{\rm ed}\right)_{\rm meas.}$	$\left(S_{ m JJ'}^{ m ed} ight)_{ m calc.}$	$\left(S_{\rm JJ'}^{\rm ed}\right)_{\rm meas.}$	$\left(S_{\rm JJ'}^{\rm ed} ight)_{ m calc.}$	
from ${}^{4}I_{15/2}$	$\times 10^{-20} {\rm cm}^2$	$\times 10^{-20} \mathrm{cm}^2$	$\times 10^{-20} {\rm cm}^2$	$\times 10^{-20} {\rm cm}^2$	$\times 10^{-20} \mathrm{cm}^2$	$\times 10^{-20} \mathrm{cm}^2$	
$^{4}I_{11/2}$	1.19484	1.02122	1.11966	0.92318	1.32161	1.06635	
$^{4}I_{9/2}$	0.37849	0.44941	0.34634	0.42192	0.41429	0.47651	
${}^{4}F_{9/2}$	2.44823	2.45268	2.24506	2.26101	2.59892	2.59826	
${}^{4}S_{3/2}$	0.47097	0.54517	0.35545	0.49048	0.48819	0.57685	
$^{2}H(2)_{11/2}$	2.30644	2.39657	2.21259	2.30857	2.21233	2.17921	
${}^{4}F_{7/2}$	1.73245	1.90524	1.56423	1.72902	1.79117	2.01678	
${}^{4}F_{5/2}$	0.87597	0.55059	0.80683	0.49536	0.91134	0.58259	
$^{2}G(1)_{9/2}$	0.43984	0.60624	0.39568	0.54412	0.46899	0.63777	
$^{4}G_{11/2}$	3.13488	3.07027	3.02660	2.95814	2.75558	2.78912	
$^{4}G_{9/2}$	1.08832	0.89686	1.02737	0.83142	1.16015	0.95036	

Table 2: Measured and calculated absorption line strengths of Er^{3+} *in* CaF_2 *.*

Table 3: J-O parameters for various ErF₃ concentrations in CaF₂ crystals.

Concentration of Er ³⁺	$\Omega_2 \times 10^{-20} \mathrm{cm}^2$	$\Omega_4 \times 10^{-20} \mathrm{cm}^2$	$\Omega_6 \times 10^{-20} \mathrm{cm}^2$	
0.15 mol%	1.62356	2.45380	2.46572	
0.30 mol%	1.61584	2.30920	2.21837	
0.65 mol%	1.21408	2.60210	2.60899	

The radiative lifetime, $\tau_{\rm rad}$, for an excited state is:

$$\tau_{\rm rad.} = \frac{1}{\sum_{\rm r} A_{\rm JJ'}} \tag{5}$$

where the sum is taken over all final lower-lying states J'. The fluorescence branching ratios, β_{JJ} , can be determined from the radiative decay rates by the following expression:

$$\beta_{\rm JJ'} = \frac{A_{\rm JJ'}}{\sum_{\rm J'} A_{\rm JJ'}} = A_{\rm JJ'} \tau_{\rm rad.}$$
(6)

All these values are given in Table 4.

Transition	0.15 mol%			0.30 mol%			0.67 mol%		
from ${}^{4}I_{15/2}$	$\begin{array}{c} A_{JJ'} \\ (s^{-1}) \end{array}$	$eta_{ m JJ'}$	$ au_{ m rad}$ (ms)	A_{JJ} , (s ⁻¹)	$eta_{ m JJ'}$	$ au_{ m rad}$ (ms)	A_{JJ} , (s ⁻¹)	$eta_{ m JJ'}$	τ _{rad} (ms)
${}^{4}I_{11/2}$	171.294	0.855	4.99	154.849	0.852	5.506	178.868	0.856	4.786
$^{4}I_{9/2}$	162.205	0.713	4.398	152.282	0.721	4.738	171.987	0.714	4.151
${}^{4}F_{9/2}$	1650.30	0.913	0.553	1521.342	0.913	0.600	1748.262	0.914	0.523
${}^{4}S_{3/2}$	1600.34	0.665	0.415	1439.810	0.664	0.461	1693.343	0.665	0.392
$^{2}H(2)_{11/2}$	2720.14	0.906	0.333	2620.266	0.908	0.346	2473.441	0.898	0.363
${}^{4}F_{7/2}$	3935.98	0.805 3	0.204	3571.948	0.801 0	0.2242 5	4166.434	0.8055	0.1933
${}^{4}F_{5/2}$	1969.05	0.510 4	0.259 2	10629.18 5	0.505 9	0.0476 0	12500.84	0.5104	0.0408
$^{2}G(1)_{9/2}$	1735.05	0.510 5	0.294 2	1566.554	0.506 0	0.3230	1822.406	0.5170	0.2837

Table 4: Predicted probability of spontaneous emission, branching ratios and radiativelifetimes of $CaF_2:Er^{+3}$.

4. Conclusions

A spectroscopic Judd-Ofelt study has been performed on ErF_3 -doped CaF_2 crystals in order to investigate its potential as a laser material. The intensity parameters Ω_t have been calculated from absorption spectra taken at room-temperature. The radiative decay rates and the radiative lifetimes for three various concentrations have been calculated.

The higest value of radiative decay rates has been obtained for 0.67 mol% ErF_3 for transition ${}^{4}\text{I}_{15/2} \rightarrow {}^{4}\text{F}_{5/2}$.

References

- [1] A. Kaminskii, Crystalline lasers, CRC Press, New York, 1996.
- [2] C. Labbe, Thesis, University of Caen, France, 2000.
- [3] S.C.Stefan, Thesis, West University of Timisoara, Romania, 2005.
- [4] B.R. Judd, Phys. Rev., 127 (3), 750 (1962).
- [5] G.S. Ofelt, J. Chem. Phys., 37 (3), 511 (1962).
- [6] D. Nicoara and I. Nicoara, Mater. Sci. Eng., A 102, L1 (1988).